

State-Dependent Orbital Modulation of X-rays in CYG X-3

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ABSTRACT

We analyze all of the available *RXTE* observations of Cyg X-3 in order to investigate the connection between the central X-ray source and its surrounding environment. The hardness–intensity diagram of Cyg X-3 displays a “shoe” shape rather than the Q-type shape commonly seen in other black hole X-ray binary, and exhibits no apparent hysteresis effect. During the γ -ray outbursts, no existing data are located in the hard and intermediate states, which suggest the absence of a significant population of non-thermal electrons when the source is in these states. For the first time, we present the orbital modulation of the X-ray light curve (LC) of all five states. The different energy band LCs are in phase with each other in all five states, and the modulation amplitude of both soft and hard X-ray LCs monotonously increases with decreasing hardness from hard to soft non-thermal states. We confirm that the modulation depth decreases with increasing energy in the hard, intermediate, and very high states, as originally reported by Zdziarski et al. However, in the soft non-thermal state, the hard X-ray modulation strength significantly increases and is even larger than the soft X-ray one. Our results rule out both wind absorption and jet origins of the hard X-ray LC modulation in the soft non-thermal state, and challenge our understanding of the states of Cyg X-3.

Subject headings: accretion, accretion disks – binaries: close – X-rays: binaries – X-rays: individual (Cyg X-3) – stars: winds, outflows

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1. Introduction

More than 40 years ago, the X-ray binary (XRB) Cyg X-3 was first detected by Giacconi et al. (1967) during an early rocket flight in 1966; it was then observed by all subsequent X-ray missions. It is located in the Galactic plane with an estimated distance of $\sim 7\text{--}9$ kpc (Dickey 1983; Predehl et al. 2000; Ling et al. 2009). The infrared (IR) observations indicate a short orbital period $P_{\text{orb}} = 4.8$ hr, and the donor is a Wolf–Rayet (WR) star (van Kerkwijk et al. 1996). The dense WR stellar wind prevents the optical photons from reaching the observer. Therefore, we know practically nothing about its system inclination and mass function. It remains uncertain whether the compact object is a black hole (BH) or a neutron star (NS, Stark & Saia 2003; Vilhu et al. 2009) up to now. The presence of a BH is favored when the X-ray and radio emission and the bolometric luminosity of Cyg X-3 (Szostek & Zdziarski 2008; Zdziarski et al. 2010) are considered.

As one of the brightest XRBs, the X-ray radiation of Cyg X-3 has been studied over the past years (e.g. McCollough et al. 1999). Investigating the *Rossi X-ray Timing Explorer* (*RXTE*) All Sky Monitor (ASM) and CGRO/BATSE light curves (LCs) of Cyg X-3, Hjalmarsdotter et al. (2008) found that the distribution of its ASM count rate is bimodal, with two peaks representing the soft and hard states in Cyg X-3, respectively. However, the distribution of ASM hardness and BATSE flux did not show the bimodal behavior. Inspecting 42 *RXTE*-pointed observations, Szostek et al. (2008) further divided the spectra into five groups based on their shape, which are ordered by decreasing flux above 20 keV. Hjalmarsdotter et al. (2009) interpreted the spectra in terms of Comptonization by a hybrid electron distribution and strong Compton reflection. Although the X-ray spectra of Cyg X-3 have been intensively studied, there have not been any comprehensive studies of the LC modulation in different states as yet.

In addition to the X-ray radiation, luminous radio jets are common features for both BH and NS XRBs (Coriat et al. 2011). Since first being seen by Gregory et al. (1972), Cyg X-3 has frequently been the brightest XRB at radio wavelengths with resolved jets (Tudose et al. 2010; Mioduszewski et al. 2001; Miller-Jones et al. 2004), thus qualifying the source as a microquasar. The relationship between the radio and X-ray was first pointed out by Watanabe et al. (1994), and further studied in the literature (e.g., McCollough et al. 1999; Koljonen et al. 2010). By inspecting its radio/X-ray correlation, Cyg X-3 can be distinguished from the canonical BH transients, e.g., the strong radio flares occur only when Cyg X-3 is in its soft state. Recently, γ -ray flares, which occurred close to the radio flares, were detected in Cyg X-3 by both *Fermi*/LAT (Abdo et al. 2009) and *AGILE* (Tavani et al. 2009).

Being the first microquasar detected in the GeV band, Cyg X-3 draws much attention

and dense multi-band observations. However, a number of fundamental questions are still unanswered. In this work, our aim is to answer the following questions. (1) Does the configuration of the stellar wind change during different states? (2) Does the hard X-ray LC orbital modulation result from wind absorption? (3) Does the low energy tail of the γ -ray emission contribute to the hard X-ray band? We first analyze the archived *RXTE* data in the next section. The results and their physical implications are presented in Section 3. We also check these results with X-ray monitoring data in Section 4. Our summary, conclusion, and discussion are in Section 5.

2. Data Analysis

RXTE has three instruments — the Proportional Counter Array (PCA), the High-Energy X-Ray Timing Experiment, and the ASM. The PCA was designed to cover an energy range of 2.0-60.0 keV and consists of five, nearly identical, large-area xenon Proportional Counter Units (PCUs). The field of view is about $1^\circ \times 1^\circ$. Each PCU has an area of $\sim 1300 \text{ cm}^2$, for a total effective area of 6500 cm^2 (Jahoda et al. 1996). The PCA is now well calibrated up to about 50 keV using the latest background models¹ and the response matrices v11.7 released on 2009 August 17 (Shaposhnikov et al. 2012).²

Figure 1 is an example of Cyg X-3 PCA observation data, which were used in Corbel et al. (2012). The typical source+background count spectra (MJD 55649.8, red symbol) are well above the background count spectrum (black symbol). Only when the source is faint, as when Cyg X-3 is in the ultrasoft state (MJD 55639, blue symbol; see Section 2.2 for the definition of the source states), may the background exceed the source+background count spectrum, resulting in a negative count rate at the high energy. In our work, we focus on the PCA data and present the energy-dependent orbital modulation for different states at energies $\leq 40 \text{ keV}$.

2.1. *RXTE*/PCA Data Reduction

The number of *RXTE* pointed-observations rapidly increased from 119 to 239 after 2009 January 1, when the Cyg X-3 γ -ray emission were confirmed by *Fermi*/LAT and *AGILE*

¹http://heasarc.gsfc.nasa.gov/docs/xte/pca_bkg_epoch.html

²<http://www.universe.nasa.gov/xrays/programs/rxte/pca/doc/rmf/pcarmf-11.7/>

observations. We analyze all *RXTE* observations of Cyg X-3 with HEASOFT version 6.10.³ For the PCA, we use the Standard2 data from all layers of PCU2, which operated during all the observations. The data are filtered with the standard criteria: the earth-limb elevation angle is larger than 10° and the spacecraft pointing offset is less than 0.02° . The background files are created using the program *pcabackest* and the latest bright source background model. Since the data used here spanned three *RXTE* gain epochs (from epoch 3 to epoch 5) and the PCA gain slowly drifted over time, we also analyze the data of Crab to get the correct channel-to-energy conversions (Belloni 2010). We first extract the background-subtracted LCs of Cyg X-3 channel by channel in 16 s time bins, and then normalize raw count rates with the quasi-simultaneous observations of Crab. Finally, the flux (Crab unit) of the energy band concerned is averaged with the inverse squares of the uncertainties as weights over the different channels for different gain epochs (Table 1). In this work, the 1σ (68%) errors are given with standard error propagation unless otherwise stated.

2.2. The Hardness–Intensity Diagram

The hardness–intensity diagram (HID), in which the total count rate is plotted as a function of hardness, is particularly useful for characterizing the behavior of BH XRBs (Fender et al. 2004). In order to compare with other BH XRBs, we construct the HID in a manner similar to Fender et al. (2004) and Koljonen et al. (2010). The hardness ratio is defined as the ratio of the flux between 3–6 keV and 10–15 keV (the hardness ratio of Crab is ~ 0.53), and the intensity is calculated as the sum of the count rates in both bands (Figure 2).

Following Koljonen et al. (2010), we divide the HID into five regions (I–V), somewhat arbitrarily, and produce their average spectra. The unweighted average spectra with data during *RXTE*/PCA gain epoch 5, when more than 60% observations were performed, are shown in Figure 3. Our results are similar to Szostek et al. (2008), except that no data above 40 keV are used here. To be consistent with their work, we specify regions I–V (of Figure 2) as the hard, intermediate, very high, soft non-thermal, and ultrasoft states (Szostek et al. 2008), respectively.

³See <http://heasarc.gsfc.nasa.gov/docs/software/lheasoft/>

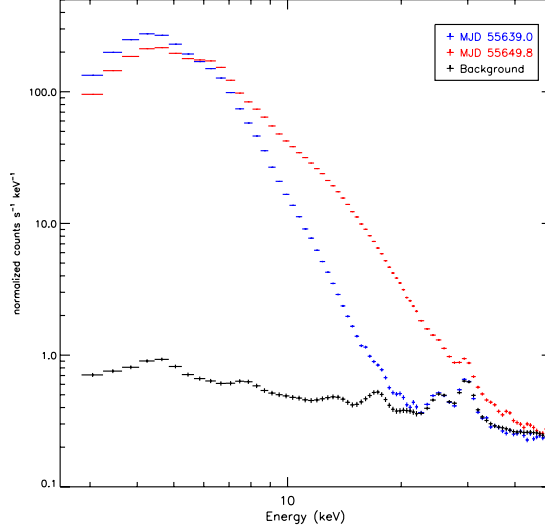


Fig. 1.— Examples of Cyg X-3 PCA count spectra. The typical source+background spectrum for the hard state (red symbol) is well above the background spectrum (black symbol). However, the background spectrum exceeded the source+background spectrum above ~ 20 keV in the ultrasoft state (blue symbol).

Table 1. Absolute Channels for Multi-band LCs in Different *RXTE* Gain Epochs

	3-6 keV	10-15 keV	2-10 keV	10-20 keV	20-40 keV	15-30 keV
epoch3	8 – 16	27 – 41	5 – 27	28 – 55	56 – 107	41 – 79
epoch4	7 – 14	24 – 35	5 – 24	25 – 47	48 – 93	35 – 69
epoch5	7 – 14	24 – 36	5 – 24	25 – 47	48 – 95	36 – 69

Note. — Channels are specified here with respect to the full 256-channel PCA. The *RXTE* epochs are documented at http://heasarc.nasa.gov/docs/xte/e-c_table.html

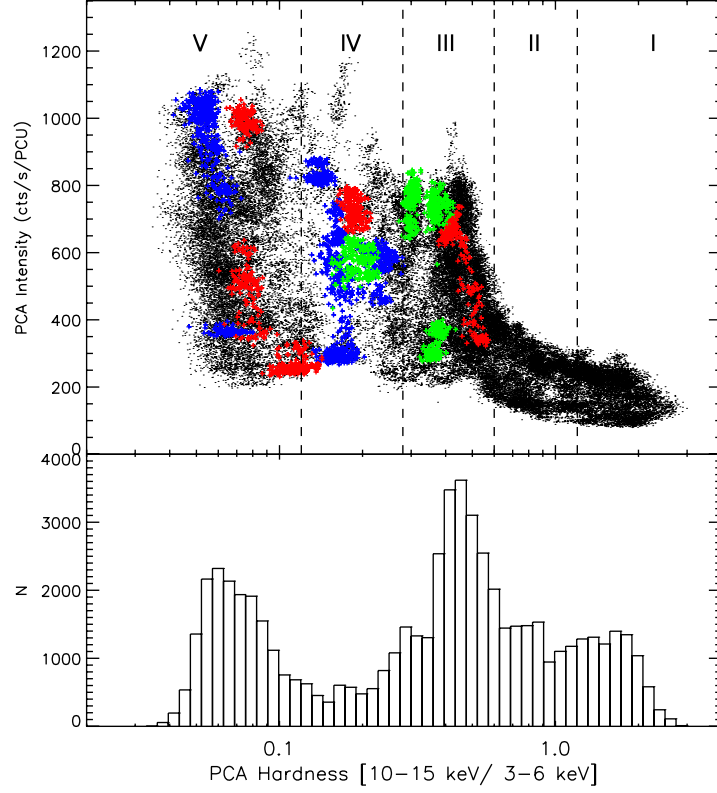


Fig. 2.— HID (upper panel) and the bimodal hardness distribution (bottom panel) of Cyg X-3, with bin size of 16 s. The HID is divided into five regions (I–V), which are occupied by the hard, intermediate, very high, soft non-thermal, and ultrasoft states, respectively. The 2009, 2010, and 2011 γ -ray outbursts data are marked with blue, green, and red crosses, respectively.

2.3. Orbital Modulation in X-Ray LC

The orbital period $P = 4.8$ hr had been detected in X-ray and IR LCs (Hanson et al. 2000; Wen et al. 2006; van Kerkwijk 1993), and recently in the GeV band (Abdo et al. 2009). Wen et al. (2006) derived a period of $P = 0.1996907(7)$ days from 8.5 yr of *RXTE*/ASM data. In addition, the evolution of the orbit had been studied in literature (Elsner et al. 1980; van der Klis & Bonnet-Bidaud 1989). The large value of the time derivative of the period (\dot{P} , $\dot{P}/P \simeq 1 \times 10^{-6} \text{yr}^{-1}$) has been suggested to be linked to the binary evolution of the system (Singh et al. 2002). Because the data used here span more than 5300 days, the derivative \dot{P} can significantly affect the orbital modulation (Figure 4). Therefore, a quadratic ephemeris is adopted for phase determination in our work.

We first calculate the phases of the data points using the ephemeris of Singh et al. (2002),

$$T_0 = 40949.392 \text{ [MJD]}, P_0 = 0.19968443 \text{ d}, \dot{P} = 5.76 \times 10^{-10}, \quad (1)$$

where P_0 is the period at the reference time T_0 . The source is variable, and the weighted average modulation profiles are biased by the higher weight data points. Alternatively, the folded LCs are given by the unweighted averages of the fluxes in each phase bin. In this work, we take two approaches to define the modulation amplitude, which are basically equivalent.

The first approach fits the folded LCs with the trigonometric function

$$f(\phi) = A * \sin[(\phi - \phi_0) * 2\pi - \pi/2] + C, \quad (2)$$

where the orbital phase ϕ is within the 0-1 interval. There are three free parameters in the function— C fluctuates near the mean value ~ 1 , ϕ_0 corresponds to the minimum flux, and A represents the modulation amplitude. However, because the LC profiles are very complex, the single sinusoid cannot accurately describe them, and may miss the actual flux extrema.

The second approach is that the modulation amplitude is calculated as $F = (f_{\max} - f_{\min}) / (f_{\max} + f_{\min})$, where f_{\max} and f_{\min} are the maximum and minimum fluxes. F is independent of the model, but might overestimate the modulation amplitude due to fluctuations of the maximum (and/or minimum) fluxes.

Because of the spectral pivoting behavior of Cyg X-3 at 10-20 keV in different states, we create the PCA 2-10 keV, 10-20 keV, and 20-40 keV LCs to investigate the energy-dependent orbital modulation. Figure 5 shows that the modulation amplitude decreases with increasing energy, $A_{2-10\text{keV}} = 0.44 \pm 0.04$, $A_{10-20\text{keV}} = 0.23 \pm 0.03$ and $A_{20-40\text{keV}} = 0.15 \pm 0.05$ ($F_{2-10\text{keV}} = 0.43 \pm 0.01$, $F_{10-20\text{keV}} = 0.31 \pm 0.01$ and $F_{20-40\text{keV}} = 0.33 \pm 0.01$) for 2-10 keV, 10-20 keV and 20-40 keV, respectively. However, the strong fluctuations exhibited in the 20-40 keV LC indicate that the modulation varies with the states. Because of the complicated

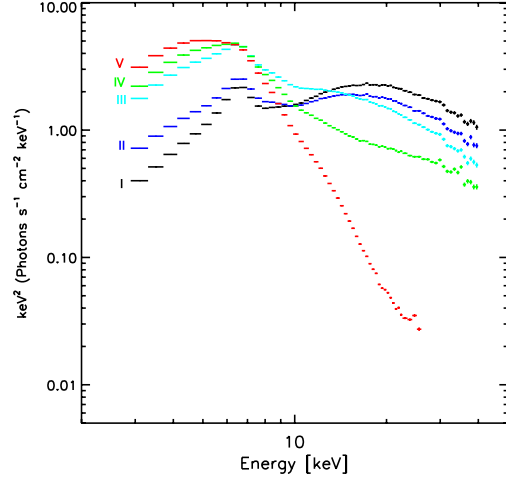


Fig. 3.— Average spectra of Cyg X-3 in different X-ray states during the *RXTE*/PCA gain epoch 5. This figure is similar to Szostek et al. (2008), except that no data above 40 keV are used in our work.

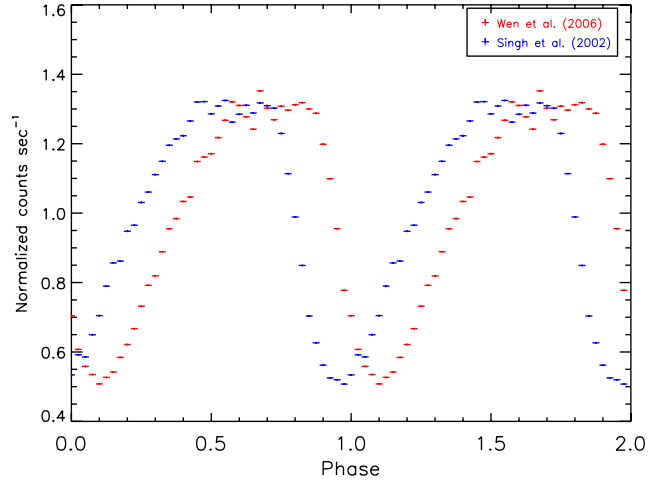


Fig. 4.— ASM dwell data over the entire lifetime of *RXTE* are folded over the orbital period using the quadratic ephemeris (blue symbol; Singh et al. 2002) and the linear ephemeris (red symbol; $T_0 = 40949.392$ [MJD], $P = 0.19969077$ d; Wen et al. 2006), respectively.

spectral states in Cyg X-3, the physical meaning of orbital modulation for all PCA data is unclear.

2.4. Energy-dependent Orbital Modulation for Different States

Because the orbital period of Cyg X-3 is unusually short, the accumulated exposure time of the *RXTE*-pointed data of ~ 980 ks covers more than 56 times its orbital period. Thus, it allows us to study the soft and hard X-ray orbital modulation in different states with a coverage time of at least ~ 5.4 times the orbital period (region IV of Figure 2) per state. The 2-10 keV, 10-20 keV, and 20-40 keV LCs are created for all five states (Figure 6), and their modulation parameters are listed in Table 2. However, due to the background being over subtracted (Figure 1), the count rate of the 20-40 keV band becomes negative in the ultrasoft state, and the modulation parameters are unavailable.

Before obtaining the final results, we carry out a series of tests.

- Boundaries between the states: the boundaries between the states are somewhat arbitrary. After shifting the boundaries by a factor of 10%, the results are approximately the same.
- Energy band of the hard X-ray: because of the background issue at 20-40 keV, the lower energy band (15-30 keV) LCs are created and used to study the hard X-ray LC modulation.
- Bin size of LCs: the short timescale variability might affect the modulation profiles and amplitude. The LCs are rebinned to 128 s (or 256 s) and then folded.

We find that all these elements and their combinations modify the modulation amplitudes only within 25%, but the modulation evolution patterns (see Section 3) of the soft and hard X-ray LCs are essentially unchanged. We fold the 256 s time-bin-size LCs and present the results in Table 3.

Since its first detection in the GeV band in 2009, *Fermi*/LAT and *AGILE* have detected several activities from Cyg X-3, which were partly covered by *RXTE* data from 2009 June 14 to 2009 July 27 (Abdo et al. 2009), 2010 May 29 to 2010 June 3 (Williams et al. 2011), and 2011 March 20 to 2011 March 24 (Corbel et al. 2012) with a total exposure time ~ 47.7 ks (2.76 orbits). The data during the 2009, 2010, and 2011 γ -ray active periods are marked by blue, green, and red crosses in the HID, respectively. The modulation profile of the LCs in 10-20 keV and 20-40 keV during these epochs deviates from the sinusoidal trace (panel VI

Table 2. Orbital Modulation for Different States

	Region I Hard State	Region II Intermediate State	Region III Very High State	Region IV Soft Non-thermal State	Region V Ultrasoft State
2-10 keV					
A	0.39 ± 0.03	0.30 ± 0.03	0.38 ± 0.03	0.41 ± 0.03	0.39 ± 0.03
ϕ_0	0.00 ± 0.01	0.96 ± 0.02	0.00 ± 0.01	0.01 ± 0.02	0.02 ± 0.01
F	0.40 ± 0.01	0.35 ± 0.01	0.41 ± 0.01	0.44 ± 0.01	0.45 ± 0.01
10-20 keV					
A	0.36 ± 0.03	0.34 ± 0.03	0.37 ± 0.03	0.48 ± 0.04	0.34 ± 0.03
ϕ_0	0.97 ± 0.01	0.95 ± 0.02	0.99 ± 0.01	0.96 ± 0.01	0.03 ± 0.01
F	0.39 ± 0.01	0.37 ± 0.01	0.40 ± 0.01	0.54 ± 0.01	0.40 ± 0.01
20-40 keV					
A	0.30 ± 0.03	0.32 ± 0.03	0.34 ± 0.03	0.57 ± 0.06	...
ϕ_0	0.95 ± 0.01	0.95 ± 0.01	0.98 ± 0.01	0.95 ± 0.02	...
F	0.35 ± 0.01	0.37 ± 0.01	0.39 ± 0.01	0.74 ± 0.03	...
15-30 keV					
A	0.32 ± 0.04	0.33 ± 0.03	0.35 ± 0.03	0.52 ± 0.04	...
ϕ_0	0.96 ± 0.01	0.95 ± 0.01	0.98 ± 0.01	0.95 ± 0.02	...
F	0.36 ± 0.01	0.37 ± 0.01	0.39 ± 0.01	0.64 ± 0.01	...

Note. — The fitted modulation parameters and their 1σ errors. Due to the background being over-subtracted, we cannot measure the modulation parameters for 20-40 keV in the ultrasoft state.

of Figure 6). Note that these γ -ray outbursts took place during different radio/X-ray states, and the accumulated exposure time (~ 47.7 ks) is relatively short. In this case, for the given phases, the LCs may be dominated by the data of different states. Thus, we further create the LCs for the very high, soft non-thermal, and ultrasoft states (see panels iii–v of Figure 6), respectively.

3. Results and Physical Implications

3.1. No Hysteresis in Cyg X-3

The HID of Cyg X-3 displays a “shoe” shape rather than the Q-type shape commonly seen in other BH XRBs (Chen et al. 2010), and exhibits no apparent hysteresis effect, which is frequently seen in other BH XRBs (Yu & Yan 2009) and a few NS XRBs (see e.g., Figure 1 in Weng & Zhang 2011). In the bottom panel of Figure 2, the PCA hardness displays a bimodal distribution, which is similar to the case in the ASM hardness of Cyg X-1 (Hjalmarsdotter et al. 2008). This is possibly due to its spectral pivot energy of ~ 10 –15 keV, which is beyond the ASM band. However, a caveat is that the PCA observations were not performed uniformly, and that the bimodal distribution might be artificial.

The strong stellar wind in Cyg X-3 provides abundant accretion material that makes the source persistent, i.e., the source never enters into the usual quiescent state. As a result, the hard state is restricted in the limited flux range. The hysteresis is thought to be triggered by the dramatic changes in accretion flow during the hydrogen ionization instability. Cyg X-3 should have a small disk due to the tight orbit and the wind accretion, and the accretion rate is higher than the critical accretion rate, making the disk temperature above the hydrogen ionization temperature everywhere (Done et al. 2007); this is probably why it does not show hysteresis. Another plausible potential candidate mechanism for suppressing the hysteresis is that the strong stellar wind collides with the accretion disk, which heats up the outer disk.

3.2. The X-Ray Behavior during the γ -Ray Outbursts

Whether the low energy tail of the γ -ray emission can give a contribution to the hard X-ray band is still under debate (Zdziarski et al. 2012b). During the γ -ray outbursts, all existing data are located in the very high, soft non-thermal, and ultrasoft states (regions III, VI, and V of Figure 2), but none in the hard and intermediate states (regions I, and II of Figure 2). This suggests the absence of a significant population of non-thermal electrons in the system when the source is in the hard and intermediate states (Hjalmarsdotter et

Table 3. Modulation Parameters for Different States

	Region I Hard State	Region II Intermediate State	Region III Very High State	Region IV Soft Non-thermal State	Region V Ultrasoft State
2-10 keV					
A	0.38 ± 0.04	0.30 ± 0.03	0.40 ± 0.03	0.43 ± 0.03	0.39 ± 0.03
ϕ_0	0.00 ± 0.01	0.96 ± 0.02	0.00 ± 0.01	0.00 ± 0.01	0.03 ± 0.01
F	0.37 ± 0.01	0.40 ± 0.01	0.40 ± 0.01	0.46 ± 0.01	0.45 ± 0.01
10-20 keV					
A	0.35 ± 0.03	0.34 ± 0.03	0.38 ± 0.03	0.50 ± 0.04	0.30 ± 0.03
ϕ_0	0.97 ± 0.01	0.95 ± 0.02	0.99 ± 0.01	0.97 ± 0.01	0.03 ± 0.01
F	0.38 ± 0.01	0.38 ± 0.01	0.39 ± 0.01	0.57 ± 0.01	0.36 ± 0.01
20-40 keV					
A	0.28 ± 0.03	0.30 ± 0.02	0.31 ± 0.02	0.56 ± 0.06	...
ϕ_0	0.95 ± 0.01	0.94 ± 0.01	0.98 ± 0.01	0.97 ± 0.01	...
F	0.33 ± 0.01	0.34 ± 0.01	0.34 ± 0.01	0.69 ± 0.04	...
15-30 keV					
A	0.31 ± 0.04	0.32 ± 0.02	0.35 ± 0.03	0.53 ± 0.05	...
ϕ_0	0.96 ± 0.01	0.95 ± 0.01	0.98 ± 0.01	0.97 ± 0.01	...
F	0.35 ± 0.01	0.35 ± 0.01	0.37 ± 0.01	0.63 ± 0.01	...

Note. — The same modulation parameters for different states as in Table 2, but now the LCs are rebinned to 256 s.

al. 2009; Acciari et al. 2011). When entering into the softer states, i.e., the very high, soft non-thermal, and ultrasoft states, the source may launch a strong radio flare. The relativistic electrons in the jets can up-scatter the dense UV photons emitted by the WR star, and produce the detected GeV emission (Dubus et al. 2010). Figure 6 shows that the modulations for the data simultaneous and not simultaneous with observations of γ -ray outbursts give the same profiles. This indicates that the X-ray and γ -ray emissions are created in different regions.

3.3. Soft and Hard X-Ray LCs Modulation

First, we construct both the soft and hard X-ray modulation profiles for all five states using all PCA data. The modulation amplitude monotonously increases from the hard to soft non-thermal states (regions I – IV). The hard X-ray has a smaller modulation depth than the soft X-ray one in the hard, intermediate, and very high states (regions I – III), which are consistent with those presented in Zdziarski et al. (2012a). However, the hard X-ray modulation strength increases significantly, and the relation between modulation amplitude and energy reverses in the soft non-thermal state (region IV).

If X-ray LC modulation is the result of the absorption in the stellar wind, the asymmetric profile, with a slow rise and a fast decline, implies a complicated wind structure. When the stellar wind becomes stronger, it provides both a denser absorbing medium and larger accretion rate, and therefore stronger modulation in softer state. At various photon energies, a given state is governed by the absorption cross sections. Both the photoionization and Compton scattering have cross sections decreasing with increasing X-ray energy (Zdziarski et al. 2012a). The hard X-ray LC shows stronger modulation than the soft X-ray LC, and therefore points to another origin of orbital modulation in the soft non-thermal state.

Both IR and γ -ray emissions were found to be also strongly modulated on the orbital period (Fender et al. 1999; Abdo et al. 2009). The IR emission was thought to arise from the WR stellar wind, which cannot emit X-ray emissions. On the other hand, the stellar radiation field is not isotropic for the relativistic electrons in the jet, and results in a modulation of the γ -ray emission (Dubus et al. 2010). In this scenario, the jet modulation profiles have opposed phases with respect to those in soft X-ray LCs, and have the maximum around the superior conjunction. However, no significant phase shift between the soft and hard X-ray LCs is detected in our data (Table 3 and Figure 6). Moreover, the modulation for the data simultaneous with the observations of γ -ray outbursts has the same features as the modulation for the other data, and thus the hard X-ray (≤ 40 keV) cannot be the low energy tail of the γ -ray emission.

4. State-Dependent Orbital Modulation from X-ray Monitoring data

The public X-ray monitoring data from the *Swift* Burst Alert Telescope (BAT) ⁴ and dwell data from *RXTE*/ASM ⁵ are also used to study the orbital modulation in different states. Cyg X-3 can transit from one state to another within one to two days, and different states may provide the same count rate, though a softer state generally emits more soft X-ray and less hard X-ray count rates (bottom panel of Figure 7; see also Figure 4 of Corbel et al. 2012). However, both BAT and ASM data only give the count rates in a single channel and lack spectral information, thus forbidding us from looking into the spectra in detail. We therefore cannot judge the source state without the help of the PCA data. In this case, we only use the data that are simultaneous with the PCA pointed observations, and assume that the source stays at the state given by the PCA data for ± 1 day (Figure 7). With these data selection criteria, the remaining data are relatively sparse and the modulation profiles might be governed by the aperiodic variability. Figure 8 shows the ASM and BAT LCs folded over its orbital period for all five states; the modulation parameters are listed in Table 4. We find that the ASM LCs have a level of modulation similar to the PCA 2–10 keV LCs, but the modulation amplitudes of BAT LCs are systematically lower than the PCA 15–30 keV LCs. Note that the BAT data encounter the same background issue as the PCA data in the ultrasoft state—the count rate becomes negative (red points in the middle panel of Figure 7), and thus no modulation parameter is available.

To check the extrema in the soft non-thermal state, we perform two more tests:

1. Because of poor statistics, the LCs are folded with a small phase bin number ($k = 10$) per period (Figure 8). Because using too large phase bin may suppress the orbital modulation, we fold the ASM and BAT LCs with $k = 15, 20$, and 30 , and obtain the modulation amplitudes $F_{\text{ASM}} = 0.49 \pm 0.01, 0.51 \pm 0.01, 0.55 \pm 0.02$, and $F_{\text{BAT}} = 0.60 \pm 0.06, 0.63 \pm 0.07, 0.67 \pm 0.08$ in the soft non-thermal state, respectively.

2. Assuming that the source stays at the state given by the PCA data for ± 0.5 and ± 2.0 days, the modulation amplitudes for the ASM and BAT data are $F_{\text{ASM}} = 0.48 \pm 0.02, 0.47 \pm 0.01$, and $F_{\text{BAT}} = 0.56 \pm 0.07, 0.52 \pm 0.04$ ($k = 15$ adopted here), respectively.

Despite systematic bias and higher statistical fluctuations, the monitoring data exhibit an evolutionary pattern similar to PCA data during these states. From the hard to soft non-thermal states, the modulation amplitude of ASM data F_{ASM} increases from 0.38 ± 0.02

⁴<http://swift.gsfc.nasa.gov/docs/swift/results/transients>

⁵http://xte.mit.edu/ASM_lc.html

to 0.49 ± 0.01 , whereas F_{BAT} increases much more, from 0.24 ± 0.03 to 0.49 ± 0.05 . The BAT LC exhibits higher (or at least the same) orbital modulation amplitude than the ASM LC in the soft non-thermal state. This implies that orbital modulation of the hard X-ray emission in the soft non-thermal state has a different origin from wind absorption.

5. Summary, Conclusion and Discussion

In this paper, we analyzed the orbital modulations of both the soft and hard X-ray LCs of Cyg X-3, and their relation with other observation data, including the GeV observations. Our main results in this work are summarized as follows:

1. The HID of Cyg X-3 displays a “shoe” shape rather than the Q-type shape commonly seen in other BH XRBs, and exhibits no apparent hysteresis effect, which is consistent with previous work (Hjalmarsdotter et al. 2008; Koljonen et al. 2010).
2. Owing to a lack of a significant population of non-thermal electrons in the system, there is no γ -ray outburst detected in the hard and intermediate states.
3. We confirm that both the soft and hard LCs have stronger orbital modulation in softer states, and that the modulation amplitude decreases with increasing energy in the hard, intermediate, and very high states; this was first reported by Zdziarski et al. (2012a). However, in the soft non-thermal state, the modulation amplitude of the hard X-ray LC becomes larger than that of the soft X-ray, which is in contradiction with the wind absorption scenario.

There is no clear evidence of jet contribution to the observed X-ray radiation, since the modulation for the data that are simultaneous with the observations of γ -ray outbursts has the same features as the modulation for the other data. The result indicates that the X-ray and γ -ray emissions are created in different regions. Because the hard X-ray LCs are always in phase with the soft X-ray LCs, but out of phase with γ -ray emission, we can further rule out the jet origin of hard X-ray and its orbital modulation.

Our results of state-dependent orbital modulation present a challenge for us in understanding the state of Cyg X-3.

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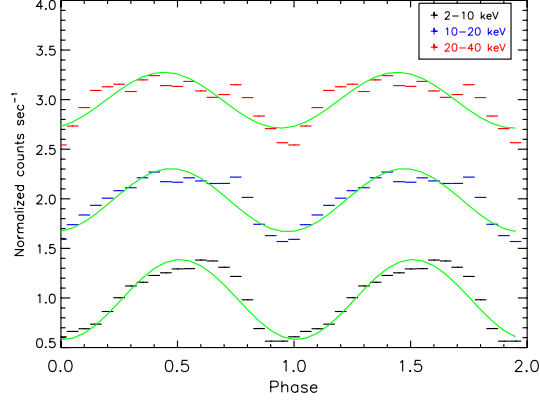


Fig. 5.— PCA LCs of 2-10 keV (black points), 10-20 keV (blue points), and 20-40 keV (red points) folded over its orbital period, respectively. The 10-20 keV and 20-40 keV LCs are shifted up by 1 and 2, respectively, for clarity. The LCs are also fitted with the trigonometric function (green lines).

Table 4. Orbital Modulation in ASM and BAT Data

	Region I Hard State	Region II Intermediate State	Region III Very High State	Region IV Soft Non-thermal State	Region V Ultrasoft State
<i>RXTE</i> /ASM 1.5-12 keV					
A	0.32 ± 0.06	0.38 ± 0.09	0.38 ± 0.05	0.43 ± 0.06	0.42 ± 0.06
ϕ_0	0.98 ± 0.03	0.00 ± 0.03	0.00 ± 0.02	0.01 ± 0.02	0.02 ± 0.03
F	0.38 ± 0.02	0.45 ± 0.01	0.39 ± 0.01	0.49 ± 0.01	0.47 ± 0.01
<i>Swift</i> /BAT 15-50 keV					
A	0.26 ± 0.04	0.28 ± 0.06	0.27 ± 0.05	0.37 ± 0.10	...
ϕ_0	0.92 ± 0.03	0.93 ± 0.03	0.96 ± 0.04	0.95 ± 0.04	...
F	0.24 ± 0.03	0.30 ± 0.02	0.34 ± 0.02	0.49 ± 0.05	...

Note. — Due to the background being over-subtracted, we cannot measure the modulation parameters for BAT LC in the ultrasoft state.

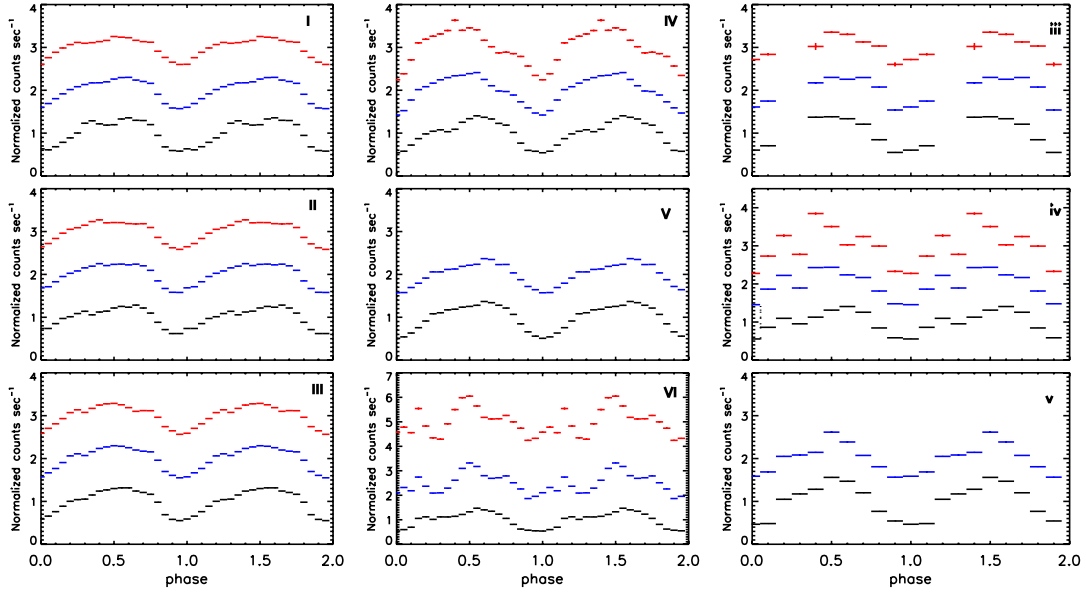


Fig. 6.— Panels I–V: the LCs of 2–10 keV (black points), 10–20 keV (blue points), and 20–40 keV (red points) folded over its orbital period for regions I–V of Figure 2, respectively. The 10–20 keV and 20–40 keV LCs in panels I, II, III, IV, V, VI, iii, iv, and v are shifted up by (1, 1, 1, 1, 1, 1.5, 1, 1, 1) and (2, 2, 2, 2, –, 4, 2, 2, –), respectively, for clarity. Panel VI: the LCs of 2–10 keV, 10–20 keV, and 20–40 keV during the 2009, 2010, and 2011 γ -ray active periods. During γ -ray outbursts, all existing data are located in the very high, soft non-thermal, and ultrasoft states, and their LCs are plotted in panels iii–v, respectively.

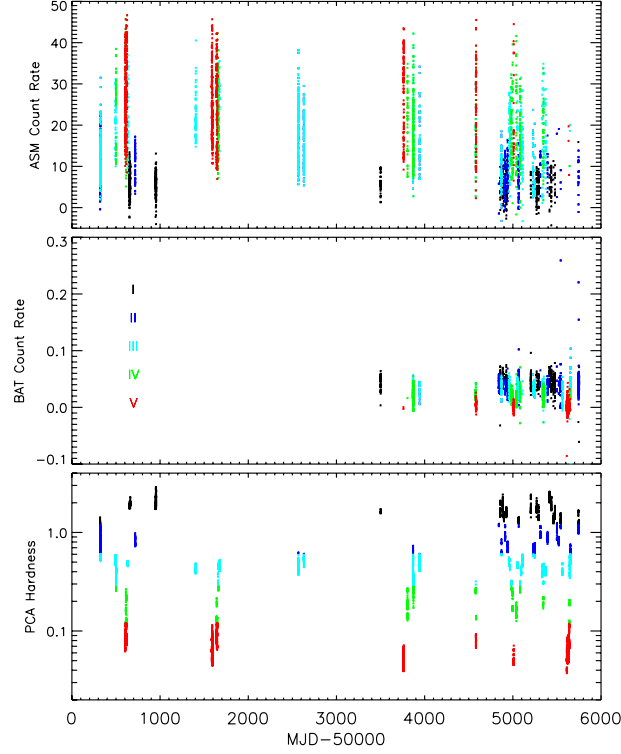


Fig. 7.— From top to bottom panels are shown the ASM (1.5-12 keV) and BAT (15-50 keV) light curves, and the temporal evolution (256 s time bin size) of the PCA hardness ratios (Figure 2), respectively. Note that we only use the ASM and BAT data which are simultaneous with the PCA-pointed observations. The data of states I, II, III, IV, and V are marked with red, green, cyan, blue and black symbols (see the text for description of the states), respectively.

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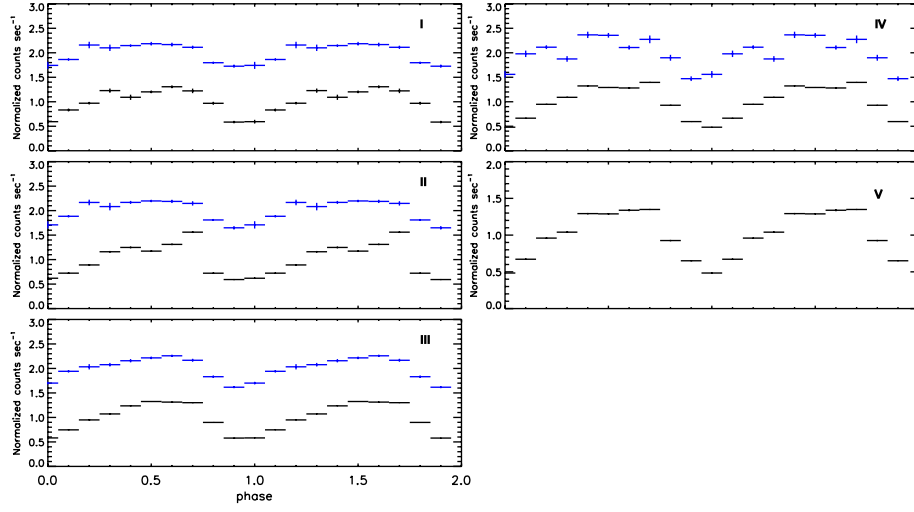


Fig. 8.— Panels I–V: The LCs of ASM (black points) and BAT (blue points) folded over its orbital period for regions I–V of Figure 2, respectively. The BAT LCs are shifted up by 1 for clarity.